

Gauge Invariant Two-Photon-Exchange Contributions in $e^-\pi^+ \rightarrow e^-\pi^+$

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Abstract

The gauge invariant two-photon exchange (TPE) contributions in $e^-\pi^+ \rightarrow e^-\pi^+$ are discussed at hadronic level. The contact term is added to keep the full amplitude gauge invariant by two methods: one is to multiply form factors with the amplitude for point-like particles and another is to construct a gauge invariant Lagrangian. The practical calculations show the TPE contributions by these two methods are almost the same, while the later method is favored when extending the discussion to processes including two charged finite-size particles like $ep \rightarrow en\pi^+$.

PACS numbers: 12.15.Lk, 13.40.Gp, 14.40.Be

Key words: Two-Photon Exchange, Form Factor, Gauge Invariant

1 Introduction

It has been shown the two-photon exchange (TPE) contributions in unpolarized elastic ep scattering play an important role in extracting the electromagnetic form factors of the proton from the angle dependence of cross section. It is natural to expect that similar effects may exist in the unpolarized $ep \rightarrow en\pi^+$, which is also used to extract the electromagnetic π form factor or σ_L from the angle dependence of cross section [1]. In the literature, many model dependent calculations [2] and model independent analyses [3] have been made to study TPE contributions in elastic ep scattering, while the TPE contributions in $ep \rightarrow en\pi^+$ are much more complex and the discussion on such TPE contributions is deficient. Formally, how to keep gauge invariance in hadronic level for such processes [4] is a non-trivial problem, since two finite-size charged particles play their roles. Before discussing the gauge invariant TPE contributions in $ep \rightarrow en\pi^+$, it is a good basis to study the gauge invariant TPE contributions in $e^-\pi^+ \rightarrow e^-\pi^+$. The TPE contributions in the latter process have been studied in [5, 6], while the contact term is usually neglected. This leads to manifest breakdown of gauge invariance. For the processes with charged non-point-like particles, the usual way to keep the full amplitude gauge invariant is to multiply form factors with the amplitudes for point-like particles. In this letter, we introduce a gauge invariant Lagrangian to treat the TPE contributions in $e^-\pi^+ \rightarrow e^-\pi^+$. Such a method can also be applied to treat the TPE contributions in $ep \rightarrow en\pi^+$ directly. We arrange our discussion as follows: in section 2, the TPE contributions in the literature are reviewed and the way to restore gauge invariance at the amplitude level is discussed; in section 3, a simple gauge invariant Lagrangian is constructed to describe the electromagnetic interactions of π and the TPE contributions are discussed by this Lagrangian; in section 4, the numerical results are presented.

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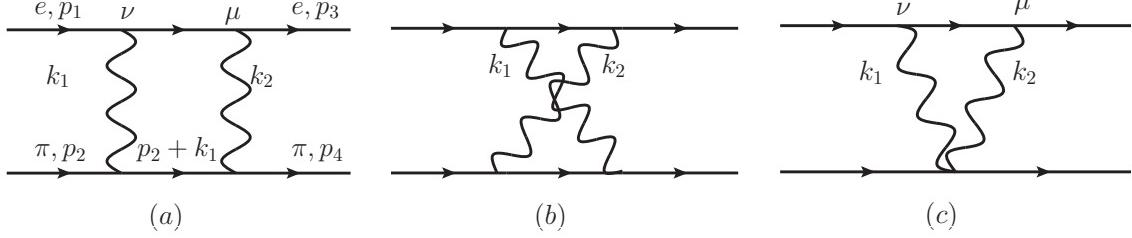


Figure 1: Two-photon-exchange diagrams with elastic intermediate state: (a) box diagram, (b) cross-box diagram and (c) contact term diagram.

2 Gauge Invariant TPE Contributions in $e^- \pi^+ \rightarrow e^- \pi^+$: A

For a charged point-like pseudoscalar particle, the electromagnetic interaction to the lowest order can be described as

$$L_0 = (D_\mu \phi)^* D^\mu \phi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}, \quad (1)$$

with $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$, $D_\mu = \partial_\mu + ie_Q A_\mu$ and e_Q being the charge of pseudoscalar particle. To keep the gauge invariance, a contact term may be introduced by the minimal coupling. This is different with point-like spin- $\frac{1}{2}$ particle where contact term is not necessary.

For finite-size charged pseudoscalar particles such as π^+ with $e_Q = -e = |e|$, higher order interactions are needed to describe its electromagnetic structure. In the literature [5, 6], to describe such structure a form factor is directly multiplied with the point-like particle vertex. This corresponds to the following replacement for the vertex:

$$ie(p_1 + p_2)_\mu \rightarrow ie(p_1 + p_2)_\mu F_\pi(q^2), \quad (2)$$

with $p_1, p_2, q \equiv p_2 - p_1$ the momentum of incoming π^+ , out coming π^+ and photon. By such replacement, it is easy to check that the sum of amplitudes corresponding to TPE diagrams Fig.1(a) and Fig.1(b) is not gauge invariant. To restore the gauge invariance, the contact term should be considered as the point-like particle case, with the simple replacement is (we named it method A):

$$i2e^2 g_{\mu\nu} \rightarrow i2e^2 g_{\mu\nu} F_\pi(k_1^2) F_\pi(k_2^2), \quad (3)$$

with k_1, k_2 the momentum of (incoming) photons and the diagram Fig.1(c) due to contact term is included. The amplitudes corresponding to the three diagrams in Fig.1 in Feynman gauge by this method read as

$$\begin{aligned} \mathcal{M}_{\gamma\gamma}^{A,(a)} &= -i \int \frac{d^4 k_1}{(2\pi)^4} \frac{\bar{u}_e(p_3)(-ie\gamma_\mu)(p_1 - k_1 + m_e)(-ie\gamma_\nu)u_e(p_1)}{[(p_1 - k_1)^2 - m_e^2 + i\epsilon][(p_2 + k_1)^2 - m_\pi^2 + i\epsilon]} \\ &\quad \frac{[ieF_\pi(k_2^2)(2p_4 - k_2)^\mu][ieF_\pi(k_1^2)(2p_2 + k_1)^\nu]}{(k_1^2 + i\epsilon)(k_2^2 + i\epsilon)}, \\ \mathcal{M}_{\gamma\gamma}^{A,(b)} &= -i \int \frac{d^4 k_1}{(2\pi)^4} \frac{\bar{u}_e(p_3)(-ie\gamma_\mu)(p_1 - k_1 + m_e)(-ie\gamma_\nu)u_e(p_1)}{[(p_1 - k_1)^2 - m_e^2 + i\epsilon][(p_2 + k_2)^2 - m_\pi^2 + i\epsilon]} \\ &\quad \frac{[ieF_\pi(k_1^2)(2p_4 - k_1)^\nu][ieF_\pi(k_2^2)(2p_2 + k_2)^\mu]}{(k_1^2 + i\epsilon)(k_2^2 + i\epsilon)}, \\ \mathcal{M}_{\gamma\gamma}^{A,(c)} &= -i \int \frac{d^4 k_1}{(2\pi)^4} \frac{\bar{u}_e(p_3)(-ie\gamma_\mu)i(p_1 - k_1 + m_e)(-ie\gamma_\nu)u_e(p_1)}{[(p_1 - k_1)^2 - m_e^2 + i\epsilon]} \\ &\quad (-) \frac{[i2e^2 F_\pi(k_2^2) F_\pi(k_1^2) g^{\mu\nu}]}{(k_1^2 + i\epsilon)(k_2^2 + i\epsilon)}, \end{aligned} \quad (4)$$

where $\mathcal{M}_{\gamma\gamma}^{A,(a)}$ and $\mathcal{M}_{\gamma\gamma}^{A,(b)}$ are the same as [5]. Now it is easy to check the full amplitude is not dependent on the gauge parameter in the photon's propagators.

With such a method based on the amplitudes directly, in principle that is not the case. And it is also not easy to extend it to processes with two finite-size charged particles in a unitary way. In the following, we construct a simple gauge invariant Lagrangian to discuss the TPE contributions.

3 Gauge Invariant TPE Contributions in $e^-\pi^+ \rightarrow e^-\pi^+$: B

Differently from using direct replacements as above, higher order terms can be added to describe the structure formally, one simple form being

$$L = L_0 + L_1, \quad (5)$$

with

$$L_1 = ie_Q D_\mu \phi^* \phi \partial_\nu f(-\partial_\rho \partial^\rho) F^{\mu\nu} + h.c.$$

Based on this Lagrangian the electromagnetic form factors of π at tree level can be written as

$$\langle p_2 | J_\mu | p_1 \rangle = (1 + q^2 f(q^2)) (p_1 + p_2)_\mu. \quad (6)$$

Comparing with the general form of electromagnetic form factor of the π

$$\langle p_2 | J_\mu | p_1 \rangle = F_\pi(q^2) (p_1 + p_2)_\mu, \quad (7)$$

the following relation is obtained

$$F_\pi(q^2) = 1 + q^2 f(q^2). \quad (8)$$

In principle, the Lagrangian Eq.(5) is not the most general one, while it is the simplest one to keep the gauge invariance in a manifest way. With Lagrangian Eq.(5), the amplitudes in Feynman gauge for the three diagrams in Fig.1 can be expressed as

$$\begin{aligned} \mathcal{M}_{\gamma\gamma}^{B,(a)} &= -i \int \frac{d^4 k_1}{(2\pi)^4} \frac{\bar{u}_e(p_3)(-ie\gamma_\mu)(p_1 - k_1 + m_e)(-ie\gamma_\nu)u_e(p_1)}{[(p_1 - k_1)^2 - m_e^2 + i\epsilon][(p_2 + k_1)^2 - m_\pi^2 + i\epsilon]} \\ &\quad \frac{\Gamma^\mu(p_4, p_4 - k_2)\Gamma^\nu(p_4 - k_2, p_2)}{(k_1^2 + i\epsilon)(k_2^2 + i\epsilon)}, \\ \mathcal{M}_{\gamma\gamma}^{B,(b)} &= -i \int \frac{d^4 k_1}{(2\pi)^4} \frac{\bar{u}_e(p_3)(-ie\gamma_\mu)(p_1 - k_1 + m_e)(-ie\gamma_\nu)u_e(p_1)}{[(p_1 - k_1)^2 - m_e^2 + i\epsilon][(p_2 + k_2)^2 - m_\pi^2 + i\epsilon]} \\ &\quad \frac{\Gamma^\mu(p_4, p_4 - k_1)\Gamma^\nu(p_4 - k_1, p_2)}{(k_1^2 + i\epsilon)(k_2^2 + i\epsilon)}, \\ \mathcal{M}_{\gamma\gamma}^{B,(c)} &= -i \int \frac{d^4 k_1}{(2\pi)^4} \frac{\bar{u}_e(p_3)(-ie\gamma_\mu)i(p_1 - k_1 + m_e)(-ie\gamma_\nu)u_e(p_1)(-\Lambda^{\mu\nu}(k_1, k_2))}{[(p_1 - k_1)^2 - m_e^2 + i\epsilon](k_1^2 + i\epsilon)(k_2^2 + i\epsilon)}, \end{aligned} \quad (9)$$

with

$$\begin{aligned} \Gamma^\mu(p_f, p_i) &= ie [(1 + f(q^2)q^2)(p_f + p_i)^\mu - f(q^2)(p_f^2 - p_i^2)q^\mu], \\ \Lambda^{\mu\nu}(k_1, k_2) &= 2ie^2 [g^{\mu\nu} + f(k_1^2)(k_1^2 g^{\mu\nu} - k_1^\mu k_1^\nu) + f(k_2^2)(k_2^2 g^{\mu\nu} - k_2^\mu k_2^\nu)]. \end{aligned} \quad (10)$$

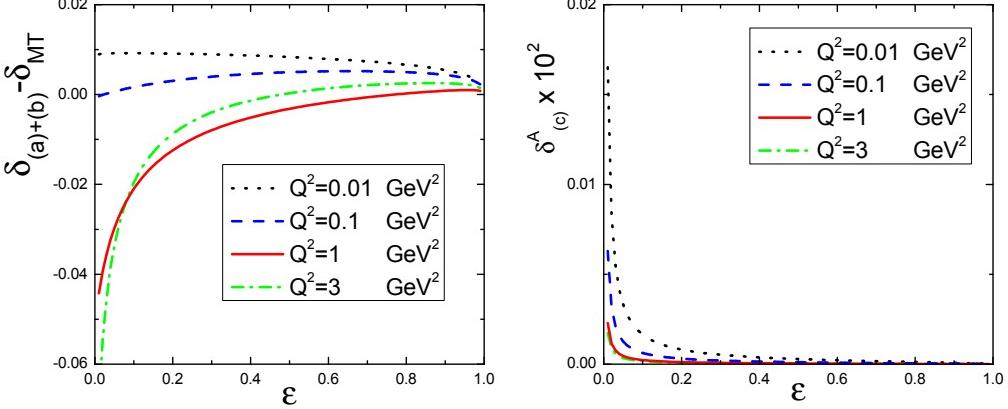


Figure 2: Two-photon-exchange contributions: the left panel is for $\delta_{(a)+(b)} - \delta_{MT}$ vs. ϵ and the right panel is for $\delta_{(c)}^A \times 10^2$ vs. ϵ both with $Q^2 = 0.01, 0.1, 1, 3 \text{ GeV}^2$.

4 Results

To show the TPE contributions, we define

$$\delta_{(a)/(b)/(c)}^{A/B} = \frac{2\text{Re}\{\mathcal{M}_0^* \mathcal{M}_{\gamma\gamma}^{A/B, (a)/(b)/(c)}\}}{|\mathcal{M}_0|^2}, \quad (11)$$

with \mathcal{M}_0 the one photon exchange amplitude, A/B refer to Method A/B and (a)/(b)/(c) refer to corresponding diagrams, respectively. In Feynman gauge, we can prove the sum $\delta_{(a)+(b)}^A$ is equal to $\delta_{(a)+(b)}^B$ with any form factors as input though $\delta_{(a)/(b)}^A$ are not equal to $\delta_{(a)/(b)}^B$, respectively. Generally such equivalence is not true for other gauge parameters. And the contributions from diagrams (c) are not equivalent by the two methods.

To show the detail, we take the same form of $F_\pi(q^2)$ with [5]

$$F_\pi(q^2) = \frac{-\Lambda^2}{q^2 - \Lambda^2}, \quad (12)$$

with $\Lambda = 0.77 \text{ GeV}$.

With this monopole form factor as input the TPE contributions can be calculated directly. And we subtract the IR divergence in the same way as [5]. The left panel of Fig.2 shows $\delta_{(a)+(b)} - \delta_{MT}$ ($\equiv \delta_{(a)+(b)}^{A/B} - \delta_{MT}$) vs. ϵ in Feynman gauge where $\epsilon = (1 + 2(1 + \tau) \tan^2(\theta/2))^{-1}$, $\tau = Q^2/4m_\pi^2$, $Q^2 = -(p_4 - p_2)^2$, θ the scattering angle and δ_{MT} denotes the correction from the box diagrams in the soft photon approximation given by the standard treatment of Mo and Tsai [7]. The right panel of Fig.2 shows $\delta_{(c)}^A$ vs. ϵ . The practical calculation shows the corrections $\delta_{(c)}^{A/B}$ in Feynman gauge are about $10^{-5} \sim 10^{-6}$ in almost all ϵ region by both two methods for Q^2 from 0.01 GeV^2 to 3 GeV^2 . The relative magnitudes $\delta_{(c)}^A / \delta_{(c)}^B$ are shown in Tab.1. An interesting property is that $\delta_{(c)}^A / \delta_{(c)}^B$ are independent on ϵ . They are very small when $Q^2 < 1 \text{ GeV}^2$ and increase with Q^2 . The small $\delta_{(c)}^{A/B}$ result in almost the same full TPE contributions by the two methods. This means the main results by [5] are kept, while this does

not mean the contact term can be neglected in other processes or other gauges. When extending the calculation to $ep \rightarrow en\pi^+$, the contributions from such term need to be considered more carefully and the method B is favored because of the manifest gauge invariance.

	$Q^2 = 0.01\text{GeV}^2$	$Q^2 = 0.1\text{GeV}^2$	$Q^2 = 1\text{GeV}^2$	$Q^2 = 0.3\text{GeV}^2$
$\delta_{(c)}^A/\delta_{(c)}^B$	1.0002	1.0014	1.0103	1.0280

Table 1: Numerical results for $\delta_{(c)}^A/\delta_{(c)}^B$ with $Q^2 = 0.01, 0.1, 1, 3\text{GeV}^2$.

5 Acknowledgment

This work is supported by the National Sciences Foundations of China under Grant No.10805009. The author gladly acknowledges the support of Theoretical Physics Center for Science Facilities for his visit where part of this work is finished.

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